# The Lévy swap market model 

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#### Abstract

Models driven by Lévy processes are attractive since they allow for better statistical fitting compared to classical diffusion models. We derive the dynamics of the forward swap rate process in a semimartingale setting and introduce a Lévy swap market model. In order to guarantee positive rates, we model the swap rates as ordinary exponentials. We start with the most distant rate which is driven by a non-homogeneous Lévy process. Via backward induction we construct the remaining swap rates such that they become martingales under the corresponding forward swap measures. Finally we show how swaptions can be priced using bilateral Laplace transforms.


## 1 Introduction

The market models of interest rate dynamics have become increasingly popular among practitioners. These models directly specify the arbitrage-free dynamics of a set of forward Libor or swap rates. The models for forward swap rates in a pure diffusion (Brownian motion) setting were developed by Jamshidian (1997), Musiela and Rutkowski (1997b) and Rutkowski (1999, 2001). More recent approaches to interest rate models involve jump-diffusions and more generally, models driven by Lévy processes. The latter are becoming increasingly popular in finance since they allow for better statistical fitting compared to classical diffusion models (see Eberlein (2001), Eberlein and Kluge (2006), Eberlein and Koval (2005)). Processes that include jumps have already been used in several papers describing the modelling of various interest rates. In the context of instantaneous, continuously compounded interest rates, Björk, Kabanov and Runggaldier (1997) extend the classical Heath, Jarrow and Morton (1992) (henceforth HJM) framework to the case of diffusion-multivariate point processes, and Björk, Di Masi, Kabanov and Runggaldier (1997) to general semimartingales. Glasserman and Kou (2003) characterized the arbitrage-free dynamics of interest rates when the term structure is modelled through forward Libor rates or forward swap rates, in the presence of both jumps and a diffusion process. They consider the case when the jumps are modelled through a finite number of marked point processes, in which case the purely discontinuous part is of bounded variation. More explicitly, they place themselves in the generalized HJM framework of Björk, Kabanov and Runggaldier (1997) and show that the simple forward rates can be embedded in an arbitrage-free model of instantaneous forward rates. Eberlein and Özkan (2005) push the forward

[^0]Libor rate approach further into a setting of more general jump processes. Apart from the HJM framework of Björk, Di Masi, Kabanov and Runggaldier (1997), they consider the Lévy setting of Eberlein and Raible (1999) and construct the discrete tenor Lévy Libor model directly through backward induction, whence extending the approach of Musiela and Rutkowski (1997a, 1997b) from the case of pure diffusions to the Lévy setting. Eberlein and Kluge (2006) price swaptions in the HJM framework using non-homogeneous Lévy processes.

In this paper, we will develop a market model of the forward swap rates by allowing the driving process to be a Lévy process. In that sense we generalize the corresponding result in Glasserman and Kou (2003). However, our approach differs in a crucial way from that of Glasserman and Kou (2003) in that we do not start by showing that the forward swap rate model can be embedded in the framework of instantaneous forward rates of Björk, Di Masi, Kabanov and Runggaldier (1997) or Eberlein and Raible (1999). In fact, as pointed out in Hunt and Kennedy (2000), the extra burden of proving that the models fall within the HJM class is unnecessary. Instead, we use a numérairebased approach and do not explicitly specify the dynamics of the instantaneous forward rates or the bond prices. The outline of such a modelling approach in a pure diffusion setting can be found in Pelsser (2000), and in Hunt and Kennedy (2000). Furthermore, we extend the backward induction method of Rutkowski $(1999,2001)$ to the case where the forward swap rates are driven by a nonhomogeneous Lévy process.

This paper is organized as follows. In Section 2 we consider the so-called regular and reverse swap market models. We extend the approach given in Hunt and Kennedy (2000), and in Pelsser (2000), from the pure diffusion setting to a more general semimartingale setting and derive the condition for those models to be arbitrage-free. In Section 3 we develop the model for forward swap rates. This setting is appropriate for practical implementations. We discuss the pricing of plain vanilla swaptions in Section 4 by using Laplace transforms.

Finally we mention that background material about Lévy processes can be found in the textbooks of Applebaum (2004), Bertoin (1996) and Sato (1999), and in a general semimartingale setting in Jacod and Shiryaev (1987). A good reference for applications of Lévy processes is Schoutens (2003).

## 2 Swap market models under a single measure

In this section we set up a market model for forward swap rates by specifying the arbitrage-free dynamics of the rates under a single measure. This is a useful approach if we want to determine the price of more sophisticated derivatives. Assume we are given a collection of dates (called tenor structure) $0<T_{0}<T_{1}<\ldots<T_{N}$, and associated to each tenor date $T_{i}$ is a zero-coupon bond that matures at this date. We denote the price at time $t$ of such a bond by $B\left(t, T_{i}\right)$.

We extend the approach given in Hunt and Kennedy (2000), and in Pelsser (2000), in order to find the arbitrage-free dynamics for the forward swap rates under the $T_{N}$-forward measure $\mathbb{P}_{N}$ (also called terminal measure ${ }^{2}$ ), and under the $T_{0}$-forward measure $\mathbb{P}_{0}$. As mentioned in Section 1 this approach is numéraire-based. We start by specifying the dynamics for the forward swap rates and then determine the necessary relationship for any corresponding term structure model to be arbitrage-free.

Note that in this section we do not explicitly assume that the driving process is a Lévy process. However, the latter can be embedded in the given setting.

We assume in the sequel that we are given a complete probability space $(\Omega, \mathcal{F}, \mathbb{P})$ equipped with a filtration $\left(\mathcal{F}_{t}\right)_{t \geq 0}$, such that the filtered probability space $\left(\Omega, \mathcal{F}, \mathbb{P},\left(\mathcal{F}_{t}\right)_{t \geq 0}\right)$ satisfies the

[^1]usual conditions, cf. Jacod and Shiryaev (1987).

### 2.1 Regular swap market model

We assume that the tenor structure $0<T_{0}<T_{1}<\ldots<T_{N}$ is given, and $\delta_{i}=T_{i}-T_{i-1}$ for $i=1, \ldots, N$, where $T_{N}$ is a fixed time horizon. $\delta_{i}$ is the fraction of a year corresponding to the period from $T_{i-1}$ to $T_{i}$. It represents the length of the $i$ th accrual period (typically a month, a quarter, or a year). We consider a family of forward swap rates $S_{i}(t):=S\left(t, T_{i}, T_{N}\right), t \in\left[0, T_{i}\right)$, which have the same maturity date $T_{N}$ for all $i=0, \ldots, N-1$. Given a family of zero-coupon bond prices $\left(B\left(t, T_{j}\right)\right)_{0 \leq t \leq T_{j}}$ for $j=0, \ldots, N$ which are assumed to be positive semimartingales, but whose dynamics does not have to be specified any further, the accrual factor for any individual swap rate from such a family is given by

$$
\begin{equation*}
C_{i, N}(t):=\sum_{j=i+1}^{N} \delta_{j} B\left(t, T_{j}\right), \tag{1}
\end{equation*}
$$

and the forward swap rate is defined as

$$
\begin{equation*}
S_{i}(t)=\frac{B\left(t, T_{i}\right)-B\left(t, T_{N}\right)}{C_{i, N}(t)}, \tag{2}
\end{equation*}
$$

(see e.g. Brigo and Mercurio (2001) and Pelsser (2000) for more details). Musiela and Rutkowski (1997b) define a probability measure $\widetilde{\mathbb{P}}_{i}=\widetilde{\mathbb{P}}_{T_{i}}$, which is equivalent to the historical probability measure $\mathbb{P}$, to be the forward swap measure associated with the dates $T_{i}$ and $T_{N}$. Under this measure the relative bond prices $\frac{B\left(t, T_{l}\right)}{\delta_{i} B\left(t, T_{i}\right)+\cdots+\delta_{N} B\left(t, T_{N}\right)},\left(t \in\left[0, T_{l} \wedge T_{i}\right]\right)$, are local martingales under $\widetilde{\mathbb{P}}_{i}$ for every $l=0, \ldots, N$.

By choosing the bond with the largest maturity $T_{N}$ to be a numéraire, the discounted accrual factor (w.r.t. that numéraire) reads

$$
\begin{equation*}
P_{t}^{i}:=\sum_{j=i+1}^{N} \frac{\delta_{j} B\left(t, T_{j}\right)}{B\left(t, T_{N}\right)}=\frac{C_{i, N}(t)}{B\left(t, T_{N}\right)}, \tag{3}
\end{equation*}
$$

which should become a local martingale under the forward measure $\mathbb{P}_{N}$. We also define for $0 \leq$ $i \leq N-1$ the following product

$$
\begin{equation*}
\Psi_{t}^{i}:=\prod_{j=0}^{i}\left(1+\delta_{j+1} S_{j+1}(t)\right) \tag{4}
\end{equation*}
$$

We follow the convention that empty sums and products are zero and one, respectively. Note that $P_{t}^{N} \equiv S_{N}(t) \equiv 0$.
We can rewrite (3) as follows by using (2),

$$
\begin{aligned}
\frac{C_{i, N}(t)}{B\left(t, T_{N}\right)} & =\frac{\delta_{i+1} B\left(t, T_{i+1}\right)+\left(\delta_{i+2} B\left(t, T_{i+2}\right)+\cdots+\delta_{N} B\left(t, T_{N}\right)\right)}{B\left(t, T_{N}\right)} \\
& =\delta_{i+1}+\delta_{i+1}\left(\frac{B\left(t, T_{i+1}\right)-B\left(t, T_{N}\right)}{C_{i+1, N}(t)} \cdot \frac{C_{i+1, N}(t)}{B\left(t, T_{N}\right)}\right)+\frac{C_{i+1, N}(t)}{B\left(t, T_{N}\right)} \\
& =\delta_{i+1}+\left(1+\delta_{i+1} S_{i+1}(t)\right) \frac{C_{i+1, N}(t)}{B\left(t, T_{N}\right)} .
\end{aligned}
$$

By substituting the $P_{t}^{i}$ from (3) into the above equation, we obtain for $i=0, \ldots, N-1$ the recursive relation

$$
\begin{equation*}
P_{t}^{i}=\delta_{i+1}+\left(1+\delta_{i+1} S_{i+1}(t)\right) \cdot P_{t}^{i+1} . \tag{5}
\end{equation*}
$$

Multiplying both sides of the equation (5) by $\Psi_{t}^{i-1}$ we obtain

$$
\Psi_{t}^{i-1} P_{t}^{i}=\delta_{i+1} \Psi_{t}^{i-1}+\Psi_{t}^{i} P_{t}^{i+1} .
$$

By backward induction, down from $i=N-1$, it follows that

$$
\begin{equation*}
P_{t}^{i}=\frac{\sum_{j=i}^{N-1} \delta_{j+1} \Psi_{t}^{j-1}}{\Psi_{t}^{i-1}}=\sum_{j=i}^{N-1} \delta_{j+1} \prod_{k=i+1}^{j}\left(1+\delta_{k} S_{k}(t)\right) \tag{6}
\end{equation*}
$$

Remark 2.1 The $P_{t}^{i}$ in (6) corresponds to the notation $s_{i i}(t)$ in Jamshidian (1997). More generally, Jamshidian (1997) defines a sum

$$
s_{i j}(t)=\sum_{k=j}^{N-1} \delta_{k+1} \prod_{l=i+1}^{k}\left(1+\delta_{l} S_{l}(t)\right)
$$

which is equivalent to $\frac{\Psi_{t}^{j-1}}{\Psi_{t}^{i-1}} P_{t}^{j}$.
Our aim is to develop an arbitrage-free model for the term structure of interest rates specified through forward swap rates $S_{i}(\cdot)$. The next theorem states the swap rate model under the terminal measure $\mathbb{P}_{N}$. This theorem slightly generalizes Theorem 5.1 in Glasserman and Kou (2003) and is an extension of the approach in Hunt and Kennedy (2000, Section 18.3) from the Gaussian driving process to a more general setting of jump processes.

Theorem 2.1 For each $i=0, \ldots, N-1$, let $\theta_{i}(\cdot)$ be a bounded $\mathbb{R}^{d}$-valued function and $G_{i}$ : $\mathbb{R}_{+} \times \mathbb{R}^{r} \rightarrow(-1, \infty)$ be a deterministic function in $G_{\mathrm{loc}}(\mu)^{3}$. Let $W^{N}$ be a standard Brownian motion in $\mathbb{R}^{d}$ with respect to $\mathbb{P}_{N}$, and $\mu$ the random measure of jumps of a semimartingale with the continuous compensator $\nu^{N}(d t, d x)=\lambda^{N}(t, d x) d t$. The dynamics of $S_{i}(\cdot), i=0, \ldots, N-1$, is assumed to satisfy

$$
\begin{equation*}
\frac{d S_{i}(t)}{S_{i}(t-)}=\alpha_{i}(t) d t+\theta_{i}(t) d W_{t}^{N}+\int_{\mathbb{R}^{r}} G_{i}(t, x)\left(\mu-\nu^{N}\right)(d t, d x) \tag{7}
\end{equation*}
$$

Then this model is arbitrage-free if

$$
\begin{align*}
\alpha_{i}(t) & =-\sum_{j=i+1}^{N-1} \frac{\delta_{j} \sum_{k=j}^{N-1} \delta_{k+1} \prod_{l=i+1}^{k}\left(1+\delta_{l} S_{l}(t-)\right) S_{j}(t-) \theta_{j}(t) \theta_{i}(t)}{\sum_{k=i}^{N-1} \delta_{k+1} \prod_{l=i+1}^{k}\left(1+\delta_{l} S_{l}(t-)\right) \cdot\left(1+\delta_{j} S_{j}(t-)\right)} \\
& +\int_{\mathbb{R}^{r}} G_{i}(t, x)\left[1-\frac{\sum_{j=i}^{N-1} \delta_{j+1} \prod_{k=i+1}^{j}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)}{\sum_{j=i}^{N-1} \delta_{j+1} \prod_{k=i+1}^{j}\left(1+\delta_{k} S_{k}(t-)\right)}\right] \lambda^{N}(t, d x) . \tag{8}
\end{align*}
$$

Proof. As we want the model (7) to be arbitrage-free, each of the $P_{t}^{i}, i=0, \ldots, N-1$, defined in (3), has to become a local martingale under the measure $\mathbb{P}_{N}$. This imposes a relationship between the finite variation term and the coefficients of the Gaussian and the discontinuous terms in (7),

[^2]which we now derive.
Applying Itô's formula (in its product rule's version) to (5) leads to
\[

$$
\begin{equation*}
d P_{t}^{i}=\left(1+\delta_{i+1} S_{i+1}(t-)\right) d P_{t}^{i+1}+\delta_{i+1} P_{t-}^{i+1} d S_{i+1}(t)+\delta_{i+1} d\left[S_{i+1}, P^{i+1}\right]_{t} . \tag{9}
\end{equation*}
$$

\]

The quadratic covariation term on the right-hand side of (9) is according to Jacod and Shiryaev (1987, Definition I.4.53)

$$
\begin{equation*}
\left[S_{i+1}, P^{i+1}\right]_{t}=\left\langle S_{i+1}^{c}, P^{i+1, c}\right\rangle_{t}+\sum_{0 \leq s \leq t} \Delta S_{i+1}(s) \Delta P_{s}^{i+1} \tag{10}
\end{equation*}
$$

where $\Delta S_{i+1}(t)$ and $\Delta P_{t}^{i+1}$ denote the jumps of $S_{i+1}(t)$ and $P_{t}^{i+1}$, respectively. Recall that the superscript $c$ indicates that we consider the the continuous martingale part.
These jumps at time $t$ follow from (7) and (6) by definition:

$$
\Delta S_{i+1}(t)=\int_{\mathbb{R}^{r}} S_{i+1}(t-) G_{i+1}(t, x) \mu(\{t\}, d x)
$$

and

$$
\begin{aligned}
\Delta P_{t}^{i+1}= & \int_{\mathbb{R}^{r}}\left(\sum_{j=i+1}^{N-1} \delta_{j+1} \prod_{k=i+2}^{j}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)\right. \\
& \left.-\sum_{j=i+1}^{N-1} \delta_{j+1} \prod_{k=i+2}^{j}\left(1+\delta_{k} S_{k}(t-)\right)\right) \mu(\{t\}, d x) .
\end{aligned}
$$

We can now express (9) more explicitly as

$$
\begin{align*}
d P_{t}^{i}= & \left(1+\delta_{i+1} S_{i+1}(t-)\right) d P_{t}^{i+1}+\delta_{i+1} P_{t-}^{i+1} d S_{i+1}(t)+\delta_{i+1} d\left\langle S_{i+1}^{c}, P^{i+1, c}\right\rangle_{t} \\
& +\delta_{i+1} S_{i+1}(t-) \int_{\mathbb{R}^{r}} G_{i+1}(t, x) \times \\
& \times\left[\sum_{j=i+1}^{N-1} \delta_{j+1} \prod_{k=i+2}^{j}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)-P_{t-}^{i+1}\right] \mu(d t, d x) \\
= & \left(1+\delta_{i+1} S_{i+1}(t-)\right) d P_{t}^{i+1}+\delta_{i+1} P_{t-}^{i+1} d S_{i+1}(t)+\delta_{i+1} d\left\langle S_{i+1}^{c}, P^{i+1, c}\right\rangle_{t} \\
& +\delta_{i+1} S_{i+1}(t-) \int_{\mathbb{R}^{r}} G_{i+1}(t, x) \times \\
& \times\left[\sum_{j=i+1}^{N-1} \delta_{j+1} \prod_{k=i+2}^{j}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)-P_{t-}^{i+1}\right]\left(\mu-\nu^{N}\right)(d t, d x) \\
& +\delta_{i+1} S_{i+1}(t-) \int_{\mathbb{R}^{r}} G_{i+1}(t, x) \times \\
& \times\left[\sum_{j=i+1}^{N-1} \delta_{j+1} \prod_{k=i+2}^{j}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)-P_{t-}^{i+1}\right] \nu^{N}(d t, d x) . \tag{11}
\end{align*}
$$

Recall that $P_{t}^{i}$ has to become a local martingale under the measure $\mathbb{P}_{N}$. Eliminating the finite variation terms - which will be set to equal zero - in the SDE (11) yields when taking (7) into
account

$$
\begin{align*}
d P_{t}^{i}= & \left(1+\delta_{i+1} S_{i+1}(t-)\right) d P_{t}^{i+1}+P_{t-}^{i+1} \delta_{i+1} S_{i+1}(t-) \theta_{i+1}(t) d W_{t}^{N} \\
& +\delta_{i+1} P_{t-}^{i+1} S_{i+1}(t-) \int_{\mathbb{R}^{r}} G_{i+1}(t, x)\left(\mu-\nu^{N}\right)(d t, d x) \\
& +\delta_{i+1} S_{i+1}(t-) \int_{\mathbb{R}^{r}} G_{i+1}(t, x) \times \\
& \times\left[\sum_{j=i+1}^{N-1} \delta_{j+1} \prod_{k=i+2}^{j}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)-P_{t-}^{i+1}\right]\left(\mu-\nu^{N}\right)(d t, d x) \tag{12}
\end{align*}
$$

Multiplying both sides of the equation (12) by $\Psi_{t-}^{i-1}$ we get

$$
\begin{align*}
& \Psi_{t-}^{i-1} d P_{t}^{i} \\
& =\Psi_{t-}^{i} d P_{t}^{i+1}+\Psi_{t-}^{i} P_{t-}^{i+1}\left(\frac{\delta_{i+1} S_{i+1}(t-)}{1+\delta_{i+1} S_{i+1}(t-)}\right) \theta_{i+1}(t) d W_{t}^{N} \\
& +\Psi_{t-}^{i} P_{t-}^{i+1}\left(\frac{\delta_{i+1} S_{i+1}(t-)}{1+\delta_{i+1} S_{i+1}(t-)}\right) \int_{\mathbb{R}^{r}} G_{i+1}(t, x)\left(\mu-\nu^{N}\right)(d t, d x) \\
& +\Psi_{t-}^{i}\left(\frac{\delta_{i+1} S_{i+1}(t-)}{1+\delta_{i+1} S_{i+1}(t-)}\right) \int_{\mathbb{R}^{r}} G_{i+1}(t, x) \times \\
& \quad \times\left[\sum_{j=i+1}^{N-1} \delta_{j+1} \prod_{k=i+2}^{j}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)-P_{t-}^{i+1}\right]\left(\mu-\nu^{N}\right)(d t, d x) \tag{13}
\end{align*}
$$

In order to obtain a non-recursive expression for $d P_{t}^{i}$ we proceed by backward induction, down from $i=N-1$. For clarity, we consider a few steps of this induction procedure.
(1) In case when $i=N-1, \Psi_{t-}^{N-2} d P_{t}^{N-1}=0$ since $P_{t}^{N}=0$ and $S_{N}=0$.
(2) If $i=N-2$ we obtain

$$
\begin{aligned}
& \Psi_{t-}^{N-3} d P_{t}^{N-2} \\
& =\Psi_{t-}^{N-2} d P_{t}^{N-1}+\Psi_{t-}^{N-2} P_{t-}^{N-1}\left(\frac{\delta_{N-1} S_{N-1}(t-)}{1+\delta_{N-1} S_{N-1}(t-)}\right) \theta_{N-1}(t) d W_{t}^{N} \\
& +\Psi_{t-}^{N-2} P_{t-}^{N-1}\left(\frac{\delta_{N-1} S_{N-1}(t-)}{1+\delta_{N-1} S_{N-1}(t-)}\right) \int_{\mathbb{R}^{r}} G_{N-1}(t, x)\left(\mu-\nu^{N}\right)(d t, d x) \\
& +\Psi_{t-}^{N-2}\left(\frac{\delta_{N-1} S_{N-1}(t-)}{1+\delta_{N-1} S_{N-1}(t-)}\right) \int_{\mathbb{R}^{r}} G_{N-1}(t, x) \times \\
& \quad \times\left[\sum_{j=N-1}^{N-1} \delta_{j+1} \prod_{k=N}^{j}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)-P_{t-}^{N-1}\right]\left(\mu-\nu^{N}\right)(d t, d x)
\end{aligned}
$$

where $\Psi_{t-}^{N-2} d P_{t}^{N-1}=0$ from the previous step (1). This yields

$$
\Psi_{t-}^{N-3} d P_{t}^{N-2}=\Psi_{t-}^{N-2} P_{t-}^{N-1}\left(\frac{\delta_{N-1} S_{N-1}(t-)}{1+\delta_{N-1} S_{N-1}(t-)}\right) \theta_{N-1}(t) d W_{t}^{N}
$$

$$
\begin{aligned}
+ & \Psi_{t-}^{N-2}\left(\frac{\delta_{N-1} S_{N-1}(t-)}{1+\delta_{N-1} S_{N-1}(t-)}\right) \int_{\mathbb{R}^{r}} G_{N-1}(t, x) \times \\
& \times\left[\sum_{j=N-1}^{N-1} \delta_{j+1} \prod_{k=N}^{j}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)\right]\left(\mu-\nu^{N}\right)(d t, d x)
\end{aligned}
$$

For the more detailed description of this induction procedure we refer to Liinev (2004).
For a general $i, i=0, \ldots, N-1$, it therefore follows that

$$
\begin{aligned}
\Psi_{t-}^{i-1} d P_{t}^{i}= & \sum_{j=i+1}^{N-1} \Psi_{t-}^{j-1} P_{t-}^{j}\left(\frac{\delta_{j} S_{j}(t-)}{1+\delta_{j} S_{j}(t-)}\right) \theta_{j}(t) d W_{t}^{N} \\
+ & \sum_{j=i+1}^{N-1} \Psi_{t-}^{j-1}\left(\frac{\delta_{j} S_{j}(t-)}{1+\delta_{j} S_{j}(t-)}\right) \int_{\mathbb{R}^{r}} G_{j}(t, x) \times \\
& \times\left[\sum_{l=j}^{N-1} \delta_{l+1} \prod_{k=j+1}^{l}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)\right]\left(\mu-\nu^{N}\right)(d t, d x),
\end{aligned}
$$

and thus

$$
\begin{align*}
d P_{t}^{i}= & P_{t-}^{i} \sum_{j=i+1}^{N-1} \frac{\Psi_{t-}^{j-1} P_{t-}^{j}}{\Psi_{t-}^{i-1} P_{t-}^{i}}\left(\frac{\delta_{j} S_{j}(t-)}{1+\delta_{j} S_{j}(t-)}\right) \theta_{j}(t) d W_{t}^{N} \\
+ & P_{t-}^{i} \sum_{j=i+1}^{N-1} \frac{\Psi_{t-}^{j-1}}{\Psi_{t-}^{i-1} P_{t-}^{i}}\left(\frac{\delta_{j} S_{j}(t-)}{1+\delta_{j} S_{j}(t-)}\right) \times \\
& \times \int_{\mathbb{R}^{r}} G_{j}(t, x)\left[\sum_{l=j}^{N-1} \delta_{l+1} \prod_{k=j+1}^{l}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)\right]\left(\mu-\nu^{N}\right)(d t, d x), \tag{14}
\end{align*}
$$

which indeed represents a local martingale. Now we can investigate the finite variation terms in (9) and (11):

$$
\begin{align*}
& \delta_{i+1} P_{t-}^{i+1} \alpha_{i+1}(t) S_{i+1}(t-) d t+\delta_{i+1} d\left\langle S_{i+1}^{c}, P^{i+1, c}\right\rangle_{t}+\delta_{i+1} S_{i+1}(t-) \times \\
& \times \int_{\mathbb{R}^{r}} G_{i+1}(t, x)\left[\sum_{j=i+1}^{N-1} \delta_{j+1} \prod_{k=i+2}^{j}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)-P_{t-}^{i+1}\right] \nu^{N}(d t, d x)=0, \tag{15}
\end{align*}
$$

where from (7) and (14)

$$
\begin{equation*}
d\left\langle S_{i+1}^{c}, P^{i+1, c}\right\rangle_{t}=S_{i+1}(t-) \theta_{i+1}(t) P_{t-}^{i+1} \sum_{j=i+2}^{N-1} \frac{\Psi_{t-}^{j-1} P_{t-}^{j}}{\Psi_{t-}^{i} P_{t-}^{i+1}}\left(\frac{\delta_{j} S_{j}(t-)}{1+\delta_{j} S_{j}(t-)}\right) \theta_{j}(t) d t \tag{16}
\end{equation*}
$$

Combining (15) and (16) we can easily find that the drift term in (7) is for $i=0, \ldots, N-1$ given
by

$$
\begin{align*}
\alpha_{i}(t) & =-\sum_{j=i+1}^{N-1} \frac{\Psi_{t-}^{j-1} P_{t-}^{j}}{\Psi_{t-}^{i-1} P_{t-}^{i}}\left(\frac{\delta_{j} S_{j}(t-)}{1+\delta_{j} S_{j}(t-)}\right) \theta_{j}(t) \theta_{i}(t) \\
& -\int_{\mathbb{R}^{r}} G_{i}(t, x)\left[\frac{\sum_{j=i}^{N-1} \delta_{j+1} \prod_{k=i+1}^{j}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)}{P_{t-}^{i}}-1\right] \lambda^{N}(t, d x) . \tag{17}
\end{align*}
$$

By taking into account the definition (4) and relation (6), we can express the right-hand side of (17) through forward swap rates and their volatilities, yielding (8).

### 2.2 Reverse swap market models

In this section we consider another type of forward swap rate model for so-called reverse swap markets. In the pure diffusion setting this model is developed in Hunt and Kennedy (2000), and Pelsser (2000). We show here that it is indeed possible to extend this model to a semimartingale setting. Hunt and Kennedy (2000) argue that the model in Section 2.1 is suitable for path-dependent products such as barrier swaptions where the pricing is done with respect to the collection of swap rates which are reset on different dates but have a common maturity date. In this section we consider the reverse situation where the family of swap rates to be modelled, has a common start date and different maturities. These are the swap rates which underlie for example spread options. Deriving the arbitrage-free dynamics for the forward swap rates in this case is analogous to the proof of Theorem 2.1 in Section 2.1. Therefore we do not go into details, but rather present the main steps for developing the relevant dynamics.

We consider again the tenor structure $0<T_{0}<T_{1}<\cdots<T_{N}$, and a series of swaps starting on date $T:=T_{0}$ having maturity dates $T_{1}, \ldots, T_{N}$, which now differ from those of Section 2.1. We define for $i=1, \ldots, N$

$$
\begin{gather*}
C_{i, N}(t):=\sum_{j=1}^{i} \delta_{j} B\left(t, T_{j}\right), \\
S_{i}(t):=S_{i}\left(t, T, T_{i}\right)=\frac{B(t, T)-B\left(t, T_{i}\right)}{C_{i, N}(t)},  \tag{18}\\
P_{t}^{i}:=\sum_{j=1}^{i} \frac{\delta_{j} B\left(t, T_{j}\right)}{B(t, T)} \tag{19}
\end{gather*}
$$

and

$$
\Pi_{t}^{i}:=\prod_{j=1}^{i}\left(1+\delta_{j} S_{j}(t)\right)
$$

Notice that $P_{t}^{i}$ in (19) should become a local martingale under the forward measure $\mathbb{P}_{T}$. Analogous to (5), it follows from (18) and (19) that, for $i=1, \ldots, N$

$$
\begin{equation*}
\left(1+\delta_{i} S_{i}(t)\right) P_{t}^{i}=\delta_{i}+P_{t}^{i-1} \tag{20}
\end{equation*}
$$

and, by multiplying both sides with $\Pi_{t}^{i-1}$,

$$
\Pi_{t}^{i} P_{t}^{i}=\delta_{i} \Pi_{t}^{i-1}+\Pi_{t}^{i-1} P_{t}^{i-1}
$$

where $P_{t}^{0} \equiv 0$ and $\Pi_{t}^{0} \equiv 1$. Using induction starting at $i=1$, it follows that

$$
\begin{equation*}
P_{t}^{i}=\frac{\sum_{j=1}^{i} \delta_{j} \Pi_{t}^{j-1}}{\Pi_{t}^{i}}=\sum_{j=1}^{i} \frac{\delta_{j}}{\prod_{k=j}^{i}\left(1+\delta_{k} S_{k}(t)\right)} \tag{21}
\end{equation*}
$$

The model we want to study has a form similar to that of (7), which means that under the measure $\mathbb{P}_{T}$ the forward swap rates $S_{i}, i=1, \ldots, N$ have the dynamics specified as in (22).

Theorem 2.2 Assume that for each $i=1, \ldots, N, \theta_{i}(\cdot)$ is a bounded $\mathbb{R}^{d}$-valued function and $G_{i}$ : $\mathbb{R}_{+} \times \mathbb{R}^{r} \rightarrow(-1, \infty)$ is a deterministic function in $G_{\mathrm{loc}}(\mu)$. Let $W^{T}$ be a standard Brownian motion in $\mathbb{R}^{d}$ with respect to the measure $\mathbb{P}_{T}$ and let $\nu^{T}(d t, d x)=\lambda^{T}(t, d x) d t$ be the $\mathbb{P}_{T}$-compensator of $\mu$. Then the model

$$
\begin{equation*}
\frac{d S_{i}(t)}{S_{i}(t-)}=\alpha_{i}(t) d t+\theta_{i}(t) d W_{t}^{T}+\int_{\mathbb{R}^{r}} G_{i}(t, x)\left(\mu-\nu^{T}\right)(d t, d x) \tag{22}
\end{equation*}
$$

is arbitrage-free if

$$
\begin{align*}
\alpha_{i}(t) & =\sum_{j=1}^{i} \frac{\Pi_{t-}^{j} P_{t-}^{j}}{\Pi_{t-}^{i} P_{t-}^{i}}\left(\frac{\delta_{j} S_{j}(t-)}{1+\delta_{j} S_{j}(t-)}\right) \theta_{j}(t) \theta_{i}(t) \\
& -\int_{\mathbb{R}^{r}} \frac{G_{i}(t, x)}{P_{t-}^{i}}\left[\sum_{j=1}^{i} \frac{\delta_{j}}{\prod_{k=j}^{i}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)}-P_{t-}^{i}\right] \lambda^{T}(t, d x) \tag{23}
\end{align*}
$$

Proof. As we want the model (22) to be arbitrage-free, each of the $P_{t}^{i}, i=1, \ldots, N$, has to be a local martingale under the measure $\mathbb{P}_{T}$. Applying Itô's product rule to (20) yields

$$
\begin{equation*}
d P_{t}^{i-1}=\left(1+\delta_{i} S_{i}(t-)\right) d P_{t}^{i}+\delta_{i} P_{t-}^{i} d S_{i}(t)+\delta_{i} d\left[S_{i}, P^{i}\right]_{t} \tag{24}
\end{equation*}
$$

By eliminating the finite variation terms in (24) we obtain

$$
\begin{align*}
d P_{t}^{i-1}= & \left(1+\delta_{i} S_{i}(t-)\right) d P_{t}^{i}+P_{t-}^{i} \delta_{i} S_{i}(t-) \theta_{i}(t) d W_{t}^{T} \\
& +\delta_{i} P_{t-}^{i} S_{i}(t-) \int_{\mathbb{R}^{r}} G_{i}(t, x)\left(\mu-\nu^{T}\right)(d t, d x)+\delta_{i} S_{i}(t-) \int_{\mathbb{R}^{r}} G_{i}(t, x) \times \\
& \times\left[\sum_{j=1}^{i} \frac{\delta_{j}}{\prod_{k=j}^{i}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)}-P_{t-}^{i}\right]\left(\mu-\nu^{T}\right)(d t, d x) \\
= & \left(1+\delta_{i} S_{i}(t-)\right) d P_{t}^{i}+P_{t-}^{i} \delta_{i} S_{i}(t-) \theta_{i}(t) d W_{t}^{T}+\delta_{i} S_{i}(t-) \int_{\mathbb{R}^{r}} G_{i}(t, x) \times \\
& \times\left[\sum_{j=1}^{i} \frac{\delta_{j}}{\prod_{k=j}^{i}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)}\right]\left(\mu-\nu^{T}\right)(d t, d x) \tag{25}
\end{align*}
$$

Multiplying (25) by $\Pi_{t-}^{i-1}$ we obtain by induction

$$
\begin{aligned}
d P_{t}^{i}= & -P_{t-}^{i} \sum_{j=1}^{i} \frac{\Pi_{t-}^{j} P_{t-}^{j}}{\Pi_{t-}^{i} P_{t-}^{i}}\left(\frac{\delta_{j} S_{j}(t-)}{1+\delta_{j} S_{j}(t-)}\right) \theta_{j}(t) d W_{t}^{T} \\
& -P_{t-}^{i} \sum_{j=1}^{i} \frac{\Pi_{t-}^{j}}{\prod_{t-}^{i} P_{t-}^{i}}\left(\frac{\delta_{j} S_{j}(t-)}{1+\delta_{j} S_{j}(t-)}\right) \times \\
& \times \int_{\mathbb{R}^{r}} G_{j}(t, x)\left[\sum_{l=1}^{j} \frac{\delta_{j}}{\prod_{k=l}^{j}\left(1+\delta_{l} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)}\right]\left(\mu-\nu^{T}\right)(d t, d x) .
\end{aligned}
$$

In order the model (22) to be arbitrage-free, we have to set the finite variation terms in (24) equal to 0 which yields

$$
\begin{aligned}
& \delta_{i} P_{t-}^{i} \alpha_{i}(t) S_{i}(t-) d t+\delta_{i} d\left\langle S_{i}^{c}, P^{i, c}\right\rangle_{t}+ \\
& +\delta_{i} S_{i}(t-) \int_{\mathbb{R}^{r}} G_{i}(t, x)\left[\sum_{j=1}^{i} \frac{\delta_{j}}{\prod_{k=j}^{i}\left(1+\delta_{k} S_{k}(t-)\left(1+G_{k}(t, x)\right)\right)}-P_{t-}^{i}\right] \nu^{T}(d t, d x)=0
\end{aligned}
$$

from where the condition (23) follows.

## 3 The Lévy swap rate model

In Section 2 we discussed a general semimartingale framework for swap rates under the terminal measure. In this section we model the swap rates under the corresponding forward swap measures. In this sense we extend the backward modelling approach of Rutkowski $(1999,2001)$ where the swap rates are driven by a Wiener process under each forward swap measure. We consider again the family of forward swap rates $S\left(t, T_{i}\right):=S\left(t, T_{i}, T_{N}\right)$ for $i=0, \ldots, N-1$, which have the same expiry date, but differ in length of the underlying swap agreement. The essence of this approach is that by fixing $T_{N}$, one constructs the model backwards in terms of maturities (thus starting from the largest maturity), specifying at each step the change of measure such that the following swap rate process is a local martingale. This type of backward modelling was also used by Eberlein and Özkan (2005) in the context of the forward Lévy Libor model.

We assume that the discrete tenor structure $0<T_{0}<T_{1}<\ldots<T_{N}$ is given, and $\delta_{i}=$ $T_{i}-T_{i-1}$ for $i=1, \ldots, N$. Since we will proceed by backward induction we set $T_{i}^{*}=T_{N-i}$ and, in particular, $T^{*}:=T_{0}^{*}=T_{N}$. Thus, we consider a "reversed" tenor structure $0<T_{N}^{*}<$ $T_{N-1}^{*}<\ldots<T_{1}^{*}<T_{0}^{*}=T_{N}$. For the explicit construction of the forward swap rates we make the following assumptions:

## Assumption 3.1

For any maturity $T_{i}, i=0, \ldots, N-1$, there exists a bounded and continuous deterministic function $\gamma\left(\cdot, T_{i}\right) \geq 0$, which represents the volatility of the forward swap rate process $S\left(., T_{i}\right)$. The bound for the volatilities will be specified later.

## Assumption 3.2

We assume that the initial term structure of interest rates, specified by bond prices $B\left(0, T_{i}\right), i=$
$0, \ldots, N$, is given and that $B\left(0, T_{i}\right)$ are strictly decreasing in the second variable, i.e. $B\left(0, T_{i}\right)>$ $B\left(0, T_{i+1}\right), i=0, \ldots, N-1$. Consequently, the initial term structure $S\left(0, T_{i}, T_{N}\right)$ of forward swap rates is given by

$$
S\left(0, T_{i}, T_{N}\right)=\frac{B\left(0, T_{i}\right)-B\left(0, T_{N}\right)}{C_{i, N}(0)}
$$

We assume a complete stochastic basis $\left(\Omega, \mathcal{F}_{T^{*}}, \mathbb{P}_{T^{*}},\left(\mathcal{F}_{t}\right)_{0<t \leq T^{*}}\right)$ to be given. Suppose again that a family of bond prices $B\left(t, T_{m}\right), m=1, \ldots, N$, is given. For any $m=1, \ldots, N-1$ the accrual factor (1) can be rewritten as

$$
\begin{equation*}
C_{N-m, N}(t)=\sum_{l=N-m+1}^{N} \delta_{l} B\left(t, T_{l}\right)=\sum_{k=0}^{m-1} \delta_{N-k} B\left(t, T_{k}^{*}\right), \quad\left(t \in\left[0, T_{N-m+1}\right]\right) . \tag{26}
\end{equation*}
$$

We also define the relative bond prices for a fixed $i=0, \ldots, N$, and for every $k=0, \ldots, N$, by

$$
\begin{equation*}
Z_{N-i+1}\left(t, T_{k}\right):=\frac{B\left(t, T_{k}\right)}{C_{i-1, N}(t)}=\frac{B\left(t, T_{k}\right)}{\delta_{i} B\left(t, T_{i}\right)+\cdots+\delta_{N} B\left(t, T_{N}\right)}, \tag{27}
\end{equation*}
$$

which can be rewritten in terms of backward dates as

$$
\begin{equation*}
Z_{m}\left(t, T_{k}^{*}\right)=\frac{B\left(t, T_{k}^{*}\right)}{C_{N-m, N}(t)}=\frac{B\left(t, T_{k}^{*}\right)}{\delta_{N-m+1} B\left(t, T_{m-1}^{*}\right)+\ldots+\delta_{N} B\left(t, T^{*}\right)}, \quad t \in\left[0, T_{k}^{*} \wedge T_{m-1}^{*}\right] \tag{28}
\end{equation*}
$$

for any fixed $m \in\{1, \ldots, N\}$. For all $t \in\left[0, T_{m}^{*}\right]$ the forward swap rate for date $T_{m}^{*}$ equals

$$
\begin{align*}
S\left(t, T_{m}^{*}\right) & =\frac{B\left(t, T_{m}^{*}\right)-B\left(t, T^{*}\right)}{\delta_{N-m+1} B\left(t, T_{m-1}^{*}\right)+\ldots+\delta_{N} B\left(t, T^{*}\right)} \\
& =Z_{m}\left(t, T_{m}^{*}\right)-Z_{m}\left(t, T^{*}\right) \tag{29}
\end{align*}
$$

Remark 3.1 Since obviously $C_{N-1, N}(t)=\delta_{N} B\left(t, T_{N}\right)=\delta_{N} B\left(t, T^{*}\right)$, it is evident that

$$
\begin{equation*}
Z_{1}\left(t, T_{k}^{*}\right)=\frac{B\left(t, T_{k}^{*}\right)}{C_{N-1, N}(t)}=\frac{B\left(t, T_{k}^{*}\right)}{\delta_{N} B\left(t, T^{*}\right)}=\frac{1}{\delta_{N}} F_{B}\left(t, T_{k}^{*}, T^{*}\right) \tag{30}
\end{equation*}
$$

where $F_{B}\left(t, T_{k}^{*}, T^{*}\right)$ stands for the forward process corresponding to the time points $T_{k}^{*}$ and $T^{*}$. The related forward martingale measure is assumed to be $\mathbb{P}_{T^{*}}$ as given above.
We postulate that the forward swap measure for the date $T^{*}$, which we now denote by $\widetilde{\mathbb{P}}_{T^{*}}$, coincides with $\mathbb{P}_{T^{*}}$. Note that in order to stress the fact that we are dealing with the forward swap measure and to avoid any confusion with a forward measure, we add the tilde.

We proceed with the backward construction of forward swap measures. We start by defining the forward swap rate for the date $T_{1}^{*}$ by postulating that $S\left(\cdot, T_{1}^{*}\right)$ is given by

$$
\begin{equation*}
S\left(t, T_{1}^{*}\right)=S\left(0, T_{1}^{*}\right) \exp \left(\int_{0}^{t} \gamma\left(s, T_{1}^{*}\right) d \widetilde{L}_{s}^{T^{*}}\right) \tag{31}
\end{equation*}
$$

where

$$
\begin{equation*}
\widetilde{L}_{t}^{T^{*}}=\int_{0}^{t} b_{s} d s+\int_{0}^{t} \sqrt{c_{s}} \widetilde{W}_{s}^{T^{*}}+\int_{0}^{t} \int_{\mathbb{R}} x\left(\mu^{L}-\widetilde{\nu}^{T^{*}, L}\right)(d s, d x) \tag{32}
\end{equation*}
$$

is a non-homogeneous Lévy process under $\widetilde{\mathbb{P}}_{T^{*}}$, i.e. $\widetilde{W}^{T^{*}}$ is a standard Brownian motion, $c_{s}$ is positive such that $\int_{0}^{T^{*}}\left(\left|b_{s}\right|+c_{s}\right) d s<\infty, \mu^{L}$ is the random measure of jumps of the process, and $\widetilde{\nu}^{T^{*}, L}(d s, d x)=F_{s}(d x) d s$ is the $\widetilde{P}_{T^{*} \text {-compensator of } \mu^{L} \text {. We assume that the Lévy measures }}$ $F_{s}$, which are measures on $\mathbb{R}$ with $F_{s}(\{0\})=0$ and $\int_{0}^{T^{*}} \int_{\mathbb{R}}\left(x^{2} \wedge 1\right) F_{s}(d x) d s<\infty$, satisfy the following additional integrability condition

$$
\int_{0}^{T^{*}} \int_{|x|>1} \exp (u x) F_{s}(d x) d s<\infty
$$

for $|u| \leq(1+\varepsilon) M$, where $M, \varepsilon$ are positive constants such that $\sum_{i=1}^{n} \gamma\left(\cdot, T_{i}^{*}\right) \leq M$. For the sake of notation we consider a 1-dimensional driving process $\left(\widetilde{L}_{t}^{T^{*}}\right)$ only. The extension to a higher dimensional process is straightforward.

The swap rate process $S\left(t, T_{1}^{*}\right)$ has to be a martingale under $\widetilde{\mathbb{P}}_{T^{*}}$. This is achieved through the specification of $\left(b_{s}\right)$. Therefore, we choose the drift characteristics $\left(b_{s}\right)$ such that

$$
\begin{align*}
& \int_{0}^{t} \gamma\left(s, T_{1}^{*}\right) b_{s} d s=-\left(\int_{0}^{t} \int_{\mathbb{R}}\left(\mathrm{e}^{\gamma\left(s, T_{1}^{*}\right) x}-1-\gamma\left(s, T_{1}^{*}\right) x\right) \widetilde{\nu}^{T^{*}, L}(d s, d x)\right. \\
&+\left.\frac{1}{2} \int_{0}^{t} c_{s} \gamma^{2}\left(s, T_{1}^{*}\right) d s\right) \tag{33}
\end{align*}
$$

With this specification $S\left(t, T_{1}^{*}\right)$ can be written as a stochastic exponential $S\left(t, T_{1}^{*}\right)=S\left(0, T_{1}^{*}\right) \mathcal{E}_{t}\left(V\left(\cdot, T_{1}^{*}\right)\right)$ where

$$
\begin{aligned}
V\left(t, T_{1}^{*}\right) & =\int_{0}^{t} \gamma\left(s, T_{1}^{*}\right) \sqrt{c_{s}} d \widetilde{W}_{s}^{T^{*}} \\
& +\int_{0}^{t} \int_{\mathbb{R}}\left(\mathrm{e}^{\gamma\left(s, T_{1}^{*}\right) x}-1\right)\left(\mu^{L}-\widetilde{\nu}^{T^{*}, L}\right)(d s, d x)
\end{aligned}
$$

The $\widetilde{\mathbb{P}}_{T^{*} \text {-dynamics of the forward swap rate is then given by }}$

$$
\begin{align*}
d S\left(t, T_{1}^{*}\right)=S\left(t-, T_{1}^{*}\right)( & \gamma\left(t, T_{1}^{*}\right) \sqrt{c_{t}} d \widetilde{W}_{t}^{T^{*}} \\
& \left.+\int_{\mathbb{R}}\left(\mathrm{e}^{\gamma\left(t, T_{1}^{*}\right) x}-1\right)\left(\mu^{L}-\widetilde{\nu}^{T^{*}, L}\right)(d t, d x)\right) \tag{34}
\end{align*}
$$

for all $t \in\left[0, T_{1}^{*}\right]$, where $\widetilde{W^{T}}=W^{T^{*}}$ and $\widetilde{\nu}^{T^{*}, L}=\nu^{T^{*}, L}$, with initial condition

$$
S\left(0, T_{1}^{*}\right)=\frac{B\left(0, T_{1}^{*}\right)-B\left(0, T^{*}\right)}{\delta_{N} B\left(0, T^{*}\right)}
$$

Now we need to specify the process $S\left(\cdot, T_{2}^{*}\right)$, and the martingale measure for the date $T_{1}^{*}$. Referring to Remark 3.1 we have that $\widetilde{\mathbb{P}}_{T^{*}}=\mathbb{P}_{T^{*}}$, and

$$
Z_{1}\left(\cdot, T_{k}^{*}\right)=\frac{1}{\delta_{N}} \frac{B\left(\cdot, T_{k}^{*}\right)}{B\left(\cdot, T^{*}\right)}=\frac{1}{\delta_{N}} F_{B}\left(\cdot, T_{k}^{*}, T^{*}\right)
$$

follows a (strictly) positive local martingale under $\widetilde{\mathbb{P}}_{T^{*}}$. The dynamics of $Z_{1}$ can be expressed in a
general form as

$$
\begin{equation*}
d Z_{1}\left(t, T_{k}^{*}\right)=Z_{1}\left(t-, T_{k}^{*}\right)\left(\alpha_{1}\left(t, T_{k}^{*}\right) d \widetilde{W}_{t}^{T^{*}}+\int_{\mathbb{R}} \beta_{1}\left(x, t, T_{k}^{*}\right)\left(\mu^{L}-\widetilde{\nu}^{T^{*}, L}\right)(d t, d x)\right) \tag{35}
\end{equation*}
$$

for some $\alpha_{1}$ and $\beta_{1}$ (index 1 refers to the corresponding index in $Z_{1}$ ), and for any date $T_{k}^{*}$. Consider now the relative bond price $Z_{2}\left(\cdot, T_{k}^{*}\right)$,

$$
\begin{equation*}
Z_{2}\left(t, T_{k}^{*}\right)=\frac{B\left(t, T_{k}^{*}\right)}{\delta_{N} B\left(t, T^{*}\right)+\delta_{N-1} B\left(t, T_{1}^{*}\right)}=\frac{Z_{1}\left(t, T_{k}^{*}\right)}{\delta_{N-1} Z_{1}\left(t, T_{1}^{*}\right)+1} . \tag{36}
\end{equation*}
$$

We further use Lemma A. 1 from the Appendix. This lemma corresponds to Lemma 2.3 in Rutkowski (2001), which we generalized to our setting. We modified it slightly to include the measure transform. Based on the dynamics (35), and by applying the lemma to the processes $G=Z_{1}\left(\cdot, T_{k}^{*}\right)$ and $H=\delta_{N-1} Z_{1}\left(\cdot, T_{1}^{*}\right)$, we define the forward swap measure associated with the date $T_{1}^{*}$, by setting its Radon-Nikodym density as the stochastic exponential at time $T_{1}^{*}$,

$$
\frac{d \widetilde{\mathbb{P}}_{T_{1}^{*}}}{d \widetilde{\mathbb{P}}_{T^{*}}}=\mathcal{E}_{T_{1}^{*}}\left(M^{1}\right),
$$

where

$$
\begin{aligned}
M_{t}^{1} & =\int_{0}^{t} \frac{\delta_{N-1} Z_{1}\left(s-, T_{1}^{*}\right) \alpha_{1}\left(s, T_{1}^{*}\right)}{1+\delta_{N-1} Z_{1}\left(s-, T_{1}^{*}\right)} d \widetilde{W}_{s}^{T^{*}} \\
& +\int_{0}^{t} \int_{\mathbb{R}}\left(\frac{\delta_{N-1} Z_{1}\left(s-, T_{1}^{*}\right) \beta_{1}\left(x, t, T_{1}^{*}\right)}{1+\delta_{N-1} Z_{1}\left(s-, T_{1}^{*}\right)}\right)\left(\mu^{L}-\widetilde{\nu}^{T^{*}, L}\right)(d s, d x) .
\end{aligned}
$$

Then

$$
\widetilde{W}_{t}^{T_{1}^{*}}=\widetilde{W}_{t}^{T^{*}}-\int_{0}^{t} \frac{\delta_{N-1} Z_{1}\left(s-, T_{1}^{*}\right) \alpha_{1}\left(s, T_{1}^{*}\right)}{1+\delta_{N-1} Z_{1}\left(s-, T_{1}^{*}\right)} d s
$$

is the forward Brownian motion for the date $T_{1}^{*}$ and

$$
\widetilde{\nu}^{T_{1}^{*}, L}=\left(1+\frac{\delta_{N-1} Z_{1}\left(s-, T_{1}^{*}\right) \beta_{1}\left(x, s, T_{1}^{*}\right)}{1+\delta_{N-1} Z_{1}\left(s-, T_{1}^{*}\right)}\right) \widetilde{\nu}^{T^{*}, L}
$$

is the $\widetilde{\mathbb{P}}_{T_{1}^{*}}$-compensator of $\mu^{L}$. According to Lemma A.1, the process $Z_{2}\left(\cdot, T_{k}^{*}\right)$ is a local martingale under the new measure $\widetilde{\mathbb{P}}_{T_{1}^{*}}$. In order to express the measure change through the swap rates instead of using the relative bond prices, we use the following relations. First notice that from (29) and (30) we obtain that

$$
\begin{equation*}
Z_{1}\left(t, T_{1}^{*}\right)=\frac{B\left(t, T_{1}^{*}\right)}{\delta_{N} B\left(t, T^{*}\right)}=S\left(t, T_{1}^{*}\right)+Z_{1}\left(t, T^{*}\right)=S\left(t, T_{1}^{*}\right)+\delta_{N}^{-1} \tag{37}
\end{equation*}
$$

Differentiating both sides of the last equality and invoking (34) and (35), we obtain

$$
\begin{align*}
& Z_{1}\left(t-, T_{1}^{*}\right) \alpha_{1}\left(t, T_{1}^{*}\right) d \widetilde{W}_{t}^{T^{*}}+Z_{1}\left(t-, T_{1}^{*}\right) \int_{\mathbb{R}} \beta_{1}\left(x, t, T_{1}^{*}\right)\left(\mu^{L}-\widetilde{\nu}^{T^{*}, L}\right)(d t, d x) \\
& =S\left(t-, T_{1}^{*}\right) \gamma_{1}\left(t, T_{1}^{*}\right) d \widetilde{W}_{t}^{T^{*}}+S\left(t-, T_{1}^{*}\right) \int_{\mathbb{R}} \gamma_{2}\left(x, t, T_{1}^{*}\right)\left(\mu^{L}-\widetilde{\nu}^{T^{*}, L}\right)(d t, d x), \tag{38}
\end{align*}
$$

where for the shorthand notation, we have set $\gamma_{1}\left(t, T_{1}^{*}\right)=\gamma\left(t, T_{1}^{*}\right) \sqrt{c_{t}}$ and $\gamma_{2}\left(x, t, T_{1}^{*}\right)=e^{\gamma\left(t, T_{1}^{*}\right) x}-$ 1. As the Gaussian and the jump part of a semimartingale do not interact (see Jacod and Shiryaev (1987, II.2.34)), in order for this equality to hold we set

$$
\begin{aligned}
& Z_{1}\left(t-, T_{1}^{*}\right) \alpha_{1}\left(t, T_{1}^{*}\right)=S\left(t-, T_{1}^{*}\right) \gamma_{1}\left(t, T_{1}^{*}\right) \\
& Z_{1}\left(t-, T_{1}^{*}\right) \beta_{1}\left(x, t, T_{1}^{*}\right)=S\left(t-, T_{1}^{*}\right) \gamma_{2}\left(x, t, T_{1}^{*}\right)
\end{aligned}
$$

Consequently, $\widetilde{W}^{T_{1}^{*}}$ is explicitly given by the formula

$$
\widetilde{W}_{t}^{T_{1}^{*}}=\widetilde{W}_{t}^{T^{*}}-\int_{0}^{t} \frac{\delta_{N-1} S\left(s-, T_{1}^{*}\right) \gamma_{1}\left(s, T_{1}^{*}\right)}{1+\delta_{N-1} \delta_{N}^{-1}+\delta_{N-1} S\left(s-, T_{1}^{*}\right)} d s
$$

and the $\widetilde{\mathbb{P}}_{T_{1}^{*-c o m p e n s a t o r ~ o f ~}} \mu^{L}$ by

$$
\widetilde{\nu}^{T_{1}^{*}, L}=\left(1+\frac{\delta_{N-1} S\left(s-, T_{1}^{*}\right) \gamma_{2}\left(x, s, T_{1}^{*}\right)}{1+\delta_{N-1} \delta_{N}^{-1}+\delta_{N-1} S\left(s-, T_{1}^{*}\right)}\right) \widetilde{\nu}^{T^{*}, L} .
$$

Now we can define the next forward swap rate $S\left(t, T_{2}^{*}\right)$ by postulating that under $\widetilde{\mathbb{P}}_{T_{1}^{*}}$

$$
\begin{equation*}
S\left(t, T_{2}^{*}\right)=S\left(0, T_{2}^{*}\right) \exp \left(\int_{0}^{t} \gamma\left(s, T_{2}^{*}\right) d \widetilde{L}_{s}^{T_{1}^{*}}\right) \tag{39}
\end{equation*}
$$

where

$$
\begin{equation*}
\widetilde{L}_{t}^{T_{1}^{*}}=\int_{0}^{t} b_{s}^{T_{1}^{*}} d s+\int_{0}^{t} \sqrt{c_{s}} \widetilde{W}_{s}^{T_{1}^{*}}+\int_{0}^{t} \int_{\mathbb{R}} x\left(\mu^{L}-\widetilde{\nu}^{T_{1}^{*}, L}\right)(d s, d x) \tag{40}
\end{equation*}
$$



$$
\begin{align*}
& \int_{0}^{t} \gamma\left(s, T_{2}^{*}\right) b_{s}^{T_{1}^{*}} d s=-\left(\int_{0}^{t} \int_{\mathbb{R}}\left(\mathrm{e}^{\gamma\left(s, T_{2}^{*}\right) x}-1-\gamma\left(s, T_{2}^{*}\right) x\right) \widetilde{\nu}^{T_{1}^{*}, L}(d s, d x)\right. \\
&\left.+\frac{1}{2} \int_{0}^{t} c_{s} \gamma^{2}\left(s, T_{2}^{*}\right) d s\right) \tag{41}
\end{align*}
$$

With this specification of the drift term, $S\left(t, T_{2}^{*}\right)$ can be written as a stochastic exponential $S\left(t, T_{2}^{*}\right)=$ $S\left(0, T_{2}^{*}\right) \mathcal{E}_{t}\left(V\left(\cdot, T_{2}^{*}\right)\right)$ where

$$
\begin{aligned}
V\left(t, T_{2}^{*}\right) & =\int_{0}^{t} \gamma\left(s, T_{2}^{*}\right) \sqrt{c_{s}} d \widetilde{W}_{s}^{T_{1}^{*}} \\
& +\int_{0}^{t} \int_{\mathbb{R}}\left(\mathrm{e}^{\gamma\left(s, T_{2}^{*}\right) x}-1\right)\left(\mu^{L}-\widetilde{\nu}^{T_{1}^{*}, L}\right)(d s, d x)
\end{aligned}
$$

The $\widetilde{\mathbb{P}}_{T_{1}^{*-d y n a m i c s ~ o f ~ t h e ~ f o r w a r d ~ s w a p ~ r a t e ~}}$ is then given by

$$
\begin{align*}
d S\left(t, T_{2}^{*}\right)=S\left(t-, T_{2}^{*}\right)( & \gamma\left(t, T_{2}^{*}\right) \sqrt{c_{t}} d \widetilde{W}_{t}^{T_{1}^{*}} \\
& \left.+\int_{\mathbb{R}}\left(\mathrm{e}^{\gamma\left(t, T_{2}^{*}\right) x}-1\right)\left(\mu^{L}-\widetilde{\nu}^{T_{1}^{*}, L}\right)(d t, d x)\right) \tag{42}
\end{align*}
$$

for all $t \in\left[0, T_{2}^{*}\right]$ with the initial condition

$$
S\left(0, T_{2}^{*}\right)=\frac{B\left(0, T_{2}^{*}\right)-B\left(0, T^{*}\right)}{\delta_{N-1} B\left(0, T_{1}^{*}\right)+\delta_{N} B\left(0, T^{*}\right)} .
$$

As the next inductive step, we need to specify the process $S\left(t, T_{3}^{*}\right)$ and the martingale measure for the date $T_{2}^{*}$. From the previous step we know that the process $Z_{2}\left(\cdot, T_{k}^{*}\right)$ follows a (strictly) positive local martingale under $\widetilde{\mathbb{P}}_{T_{1}^{*}}$. In general, the dynamics of $Z_{2}$ can then be expressed as

$$
\begin{equation*}
d Z_{2}\left(t, T_{k}^{*}\right)=Z_{2}\left(t-, T_{k}^{*}\right)\left(\alpha_{2}\left(t, T_{k}^{*}\right) d \widetilde{W}_{t}^{T_{1}^{*}}+\int_{\mathbb{R}} \beta_{2}\left(x, t, T_{k}^{*}\right)\left(\mu^{L}-\widetilde{\nu}^{T_{1}^{*}, L}\right)(d t, d x)\right) \tag{43}
\end{equation*}
$$

for some $\alpha_{2}$ and $\beta_{2}$ (index 2 now refers to the corresponding index in $Z_{2}$ ), and for any date $T_{k}^{*}$. At this step we consider the relative bond prices $Z_{3}\left(\cdot, T_{k}^{*}\right)$,

$$
\begin{equation*}
Z_{3}\left(t, T_{k}^{*}\right)=\frac{B\left(t, T_{k}^{*}\right)}{\delta_{N-2} B\left(t, T_{2}^{*}\right)+\delta_{N-1} B\left(t, T_{1}^{*}\right)+\delta_{N} B\left(t, T^{*}\right)}=\frac{Z_{2}\left(t, T_{k}^{*}\right)}{1+\delta_{N-2} Z_{2}\left(t, T_{2}^{*}\right)} . \tag{44}
\end{equation*}
$$

Applying Lemma A. 1 to processes $G=Z_{2}\left(\cdot, T_{k}^{*}\right)$ and $H=\delta_{N-2} Z_{2}\left(\cdot, T_{2}^{*}\right)$, we define the forward swap measure associated with date $T_{2}^{*}$, by setting its Radon-Nikodym density as the stochastic exponential at time $T_{2}^{*}$,

$$
\frac{d \widetilde{\mathbb{P}}_{T_{2}^{*}}}{d \widetilde{\mathbb{P}}_{T_{1}^{*}}}=\mathcal{E}_{T_{2}^{*}}\left(M^{2}\right),
$$

where

$$
\left.\begin{array}{rl}
M_{t}^{2} & =\int_{0}^{t} \frac{\delta_{N-2} Z_{2}\left(s-, T_{2}^{*}\right) \alpha_{2}\left(s, T_{2}^{*}\right)}{1+\delta_{N-2} Z_{2}\left(s-, T_{2}^{*}\right)} d \widetilde{W}_{s}^{T_{1}^{*}} \\
& +\int_{0}^{t} \int_{\mathbb{R}}\left(\frac{\delta_{N-2} Z_{2}\left(s-, T_{2}^{*}\right) \beta_{2}\left(x, s, T_{2}^{*}\right)}{1+\delta_{N-2} Z_{2}\left(s-, T_{2}^{*}\right)}\right)\left(\mu^{L}-\widetilde{\nu}_{1}^{*}, L\right.
\end{array}\right)(d s, d x) .
$$

Then

$$
\widetilde{W}_{t}^{T_{2}^{*}}=\widetilde{W}_{t}^{T_{1}^{*}}-\int_{0}^{t} \frac{\delta_{N-2} Z_{2}\left(s-, T_{2}^{*}\right) \alpha_{2}\left(s, T_{2}^{*}\right)}{1+\delta_{N-2} Z_{2}\left(s-, T_{2}^{*}\right)} d s \quad\left(t \in\left[0, T_{2}^{*}\right]\right)
$$

is the forward Brownian motion for the date $T_{2}^{*}$ and

$$
\widetilde{\nu}^{T_{2}^{*}, L}=\left(1+\frac{\delta_{N-2} Z_{2}\left(s-, T_{2}^{*}\right) \beta_{2}\left(x, s, T_{2}^{*}\right)}{1+\delta_{N-2} Z_{2}\left(s-, T_{2}^{*}\right)}\right) \widetilde{\nu}^{T_{1}^{*}, L} .
$$

is the $\widetilde{\mathbb{P}}_{T_{2}^{*}}$-compensator of $\mu^{L}$. According to Lemma A.1, the process $Z_{3}\left(\cdot, T_{k}^{*}\right)$ is a local martingale under the new measure $\widetilde{\mathbb{P}}_{T_{2}^{*}}$. In order to express the measure change through the swap rates instead
of using the relative bond prices, notice that

$$
\begin{aligned}
& Z_{2}\left(t, T_{2}^{*}\right) \stackrel{(28)}{=} \frac{B\left(t, T_{2}^{*}\right)}{\delta_{N-1} B\left(t, T_{1}^{*}\right)+\delta_{N} B\left(t, T^{*}\right)} \\
& \stackrel{(29)}{=} S\left(t, T_{2}^{*}\right)+Z_{2}\left(t, T^{*}\right) \\
& \stackrel{(36)}{=} S\left(t, T_{2}^{*}\right)+\frac{Z_{1}\left(t, T^{*}\right)}{1+\delta_{N-1} Z_{1}\left(t, T^{*}\right)+\delta_{N-1} S\left(t, T_{1}^{*}\right)}
\end{aligned}
$$

where the process $Z_{1}\left(\cdot, T^{*}\right)$ is already known from the previous step, cfr. (37). Differentiating the last equality we can thus find the coefficients of the Gaussian and the discontinuous terms of the process $Z_{2}\left(\cdot, T^{*}\right)$, and consequently define $\widetilde{\mathbb{P}}_{T_{2}^{*}}$.

Extension to the general case is straightforward. Consider the induction step with respect to $m$. Suppose that we have already defined the forward swap rates $S\left(t, T_{1}^{*}\right), \ldots, S\left(t, T_{m}^{*}\right)$, and specified the forward swap measure $\widetilde{\mathbb{P}}_{T_{m-1}^{*}}$. At this step we would like to determine the forward swap measure $\widetilde{\mathbb{P}}_{T_{m}^{*}}$, and the forward swap rate $S\left(\cdot, T_{m+1}^{*}\right)$. We consider the relative bond prices

$$
Z_{m+1}\left(t, T_{k}^{*}\right)=\frac{B\left(t, T_{k}^{*}\right)}{C_{N-(m+1), N}(t)}=\frac{B\left(t, T_{k}^{*}\right)}{\delta_{N-m} B\left(t, T_{m}^{*}\right)+\cdots+\delta_{N} B\left(t, T^{*}\right)}=\frac{Z_{m}\left(t, T_{k}^{*}\right)}{1+\delta_{N-m} Z_{m}\left(t, T_{m}^{*}\right)}
$$

Applying Lemma A. 1 to processes $G=Z_{m}\left(\cdot, T_{k}^{*}\right)$ and $H=\delta_{N-m} Z_{m}\left(\cdot, T_{m}^{*}\right)$, it is clear that we can define the forward swap measure associated with date $T_{m}^{*}$, by setting its Radon-Nikodym density as

$$
\begin{equation*}
\frac{d \widetilde{\mathbb{P}}_{T_{m}^{*}}}{d \widetilde{\mathbb{P}}_{T_{m-1}^{*}}}=\mathcal{E}_{T_{m}^{*}}\left(M^{m}\right) \tag{45}
\end{equation*}
$$

where

$$
\begin{align*}
M_{t}^{m} & =\int_{0}^{t} \frac{\delta_{N-m} Z_{m}\left(s-, T_{m}^{*}\right) \alpha_{m}\left(s, T_{m}^{*}\right)}{1+\delta_{N-m} Z_{m}\left(s-, T_{m}^{*}\right)} d \widetilde{W}_{s}^{T_{m-1}^{*}} \\
& +\int_{0}^{t} \int_{\mathbb{R}}\left(\frac{\delta_{N-m} Z_{m}\left(s-, T_{m}^{*}\right) \beta_{m}\left(x, s, T_{m}^{*}\right)}{1+\delta_{N-m} Z_{m}\left(s-, T_{m}^{*}\right)}\right)\left(\mu^{L}-\widetilde{\nu}^{T_{m-1}^{*}, L}\right)(d s, d x) \tag{46}
\end{align*}
$$

Then

$$
\widetilde{W}_{t}^{T_{m}^{*}}=\widetilde{W}_{t}^{T_{m-1}^{*}}-\int_{0}^{t} \frac{\delta_{N-m} Z_{m}\left(s-, T_{m}^{*}\right) \alpha_{m}\left(s, T_{m}^{*}\right)}{1+\delta_{N-m} Z_{m}\left(s-, T_{m}^{*}\right)} d s \quad\left(t \in\left[0, T_{m}^{*}\right]\right)
$$

is the forward Brownian motion for the date $T_{m}^{*}$ and

$$
\widetilde{\nu}^{T_{m}^{*}, L}=\left(1+\frac{\delta_{N-m} Z_{m}\left(s-, T_{m}^{*}\right) \beta_{m}\left(x, s, T_{m}^{*}\right)}{1+\delta_{N-m} Z_{m}\left(s-, T_{m}^{*}\right)}\right) \widetilde{\nu}^{T_{m-1}^{*}, L}
$$

is the $\widetilde{\mathbb{P}}_{T_{m}^{*}}$-compensator of $\mu^{L}$. Therefore it suffices to analyze the process

$$
Z_{m}\left(t, T_{m}^{*}\right)=\frac{B\left(t, T_{m}^{*}\right)}{\delta_{N-m+1} B\left(t, T_{m-1}^{*}\right)+\cdots+\delta_{N} B\left(t, T^{*}\right)}=S\left(t, T_{m}^{*}\right)+Z_{m}\left(t, T^{*}\right)
$$

where

$$
Z_{m}\left(t, T^{*}\right)=\frac{Z_{m-1}\left(t, T^{*}\right)}{1+\delta_{N-m+1} Z_{m-1}\left(t, T^{*}\right)+\delta_{N-m+1} S\left(t, T_{m-1}^{*}\right)} .
$$

The process $Z_{m-1}\left(\cdot, T^{*}\right)$ is known from the preceding step and is a function of the forward swap rates $S\left(t, T_{1}^{*}\right), \ldots, S\left(t, T_{m-1}^{*}\right)$. Consequently, we can express $Z_{m}\left(\cdot, T^{*}\right)$ using the terms $S\left(t, T_{1}^{*}\right), \ldots, S\left(t, T_{m-1}^{*}\right)$ and their volatilities. Having found the probability measure $\widetilde{\mathbb{P}}_{T_{m}^{*}}$, we introduce the forward swap rate $S\left(t, T_{m+1}^{*}\right)$,

$$
S\left(t, T_{m+1}^{*}\right)=S\left(0, T_{m+1}^{*}\right) \exp \left(\int_{0}^{t} \gamma\left(s, T_{m+1}^{*}\right) d \widetilde{L}_{s}^{T_{m}^{*}}\right)
$$

and so forth.
The precise form of $\frac{d \widetilde{\mathbb{P}}_{T_{m}^{*}}}{d \mathbb{\mathbb { P }}_{T_{m-1}^{*}}} \equiv \frac{d \widetilde{\mathbb{P}}_{T_{i}}}{d \mathbb{P}_{T_{i+1}}}$ in (45)-(46) is easy to find by using the following reasoning. Recall that $\widetilde{\mathbb{P}}_{i}$ is the probability measure associated to the numéraire $C_{i-1, N}$ and $\mathbb{P}_{N}$ is the probability measure associated to the numéraire $B\left(\cdot, T_{N}\right)$. According to the change of numéraire theorem (see Geman, El Karoui and Rochet (1995)) we set

$$
\rho_{t}^{i}:=\left.\frac{d \widetilde{\mathbb{P}}_{i}}{d \mathbb{P}_{N}}\right|_{\mathcal{F}_{t}}=\frac{C_{i-1, N}(t)}{B\left(t, T_{N}\right)} \cdot \frac{B\left(0, T_{N}\right)}{C_{i-1, N}(0)}=\frac{P_{t}^{i-1}}{P_{0}^{i-1}} .
$$

In view of Girsanov's theorem and equation (14) we can deduce the Radon-Nikodym derivative $\rho_{t}^{i}$ as a stochastic exponential given by

$$
\left.\frac{d \widetilde{\mathbb{P}}_{i}}{d \mathbb{P}_{N}}\right|_{\mathcal{F}_{t}}=\mathcal{E}_{t}\left(\int_{0} \varphi_{t}^{i} d W_{t}^{N}+\int_{0} \int_{\mathbb{R}}\left(Y^{i}(t, x)-1\right)\left(\mu^{L}-\nu^{N}\right)(d t, d x)\right)=: \mathcal{E}_{t}\left(X^{i}\right)
$$

where

$$
\begin{equation*}
\varphi_{t}^{i}=\sum_{j=i}^{N-1} \frac{\Psi_{t-}^{j-1} P_{t-}^{j}}{\Psi_{t-}^{i-2} P_{t-}^{i-1}}\left(\frac{\delta_{j} S_{j}(t-)}{1+\delta_{j} S_{j}(t-)}\right) \gamma_{1}\left(t, T_{j}\right) \tag{47}
\end{equation*}
$$

and

$$
\begin{align*}
Y^{i}(t, x)-1= & \sum_{j=i}^{N-1} \frac{\Psi_{t-}^{j-1}}{\Psi_{t-}^{i-2} P_{t-}^{i-1}}\left(\frac{\delta_{j} S_{j}(t-)}{1+\delta_{j} S_{j}(t-)}\right) \gamma_{2}\left(x, t, T_{j}\right) \times \\
& \times\left[\sum_{l=j}^{N-1} \delta_{l+1} \prod_{k=j+1}^{l}\left(1+\delta_{k} S_{k}(t-)\left(1+\gamma_{2}\left(x, t, T_{k}\right)\right)\right)\right] . \tag{48}
\end{align*}
$$

We proceed by writing

$$
\frac{d \widetilde{\mathbb{P}}_{i}}{d \widetilde{\mathbb{P}}_{k}}=\frac{d \widetilde{\mathbb{P}}_{i}}{d \mathbb{P}_{N}} \frac{d \mathbb{P}_{N}}{d \widetilde{\mathbb{P}}_{k}}=\frac{d \widetilde{\mathbb{P}}_{i}}{d \mathbb{P}_{N}} \frac{1}{\frac{d \mathbb{\mathbb { P }}_{k}}{d \mathbb{P}_{N}}}=\mathcal{E}\left(X^{i}\right) \frac{1}{\mathcal{E}\left(X^{k}\right)}
$$

By applying Lemmas 2.4 and 2.6 from Kallsen and Shiryaev (2002) and taking into account that $\mathcal{E}(U) \mathcal{E}(V)=\mathcal{E}(U+V+[U, V])$ for any two semimartingales $U$ and $V$, we finally obtain the
connection between any two forward swap measures $\widetilde{\mathbb{P}}_{i}$ and $\widetilde{\mathbb{P}}_{k}$ :

$$
\begin{equation*}
\frac{d \widetilde{\mathbb{P}}_{i}}{d \widetilde{\mathbb{P}}_{k}}=\mathcal{E}(Z) \tag{49}
\end{equation*}
$$

with

$$
\begin{align*}
Z_{t}= & \int_{0}^{t}\left(\sum_{j=i}^{N-1} \frac{\xi_{s-}^{j} \gamma_{1}\left(s, T_{j}\right)}{\Psi_{s-}^{i-2} P_{s-}^{i-1}}-\sum_{j=k}^{N-1} \frac{\xi_{s-}^{j} \gamma_{1}\left(s, T_{j}\right)}{\Psi_{s-}^{k-2} P_{s-}^{k-1}}\right) d \widetilde{W}_{s}^{k} \\
& +\int_{0}^{t} \int_{\mathbb{R}}\left(\frac{\sum_{j=i}^{N-1} \frac{\zeta_{j}(s, x)}{\Psi_{s-}^{i-2} P_{s-}^{i-1}}+1}{\sum_{j=k}^{N-1} \frac{\zeta_{j}(s, x)}{\Psi_{s-}^{k-2} P_{s-}^{k-1}}+1}-1\right)\left(\mu^{L}-\widetilde{\nu}^{k, L}\right)(d s, d x) \tag{50}
\end{align*}
$$

where $\xi_{s-}^{j}=\Psi_{s-}^{j-1} P_{s-}^{j}\left(\frac{\delta_{j} S_{j}(s-)}{1+\delta_{j} S_{j}(s-)}\right)$,
$\zeta_{j}(s, x)=\Psi_{s-}^{j-1}\left(\frac{\delta_{j} S_{j}(s-)}{1+\delta_{j} S_{j}(s-)}\right) \gamma_{2}\left(x, s, T_{j}\right)\left[\sum_{l=j}^{N-1} \delta_{l+1} \prod_{v=j+1}^{l}\left(1+\delta_{v} S_{v}(s-)\left(1+\gamma_{2}\left(x, s, T_{v}\right)\right)\right)\right]$,
$\widetilde{W}^{k}$ is the standard Brownian motion with respect to the forward swap measure $\widetilde{\mathbb{P}}_{k}$ and $\widetilde{\nu}^{k, L}$ is the $\widetilde{\mathbb{P}}_{k}$-compensator of $\mu^{L}$. By using (49)-(50) we easily obtain the connection between the two consecutive forward swap measures $\widetilde{\mathbb{P}}_{i}$ and $\widetilde{\mathbb{P}}_{i+1}$. Notice that the process driving the most distant forward swap rate $S\left(t, T_{N-1}, T_{N}\right)$ is a non-homogeneous Lévy process. This is however not true for the process driving $S\left(t, T_{i}, T_{N}\right), i<N-1$, constructed during the backward induction, since the associated compensators $\widetilde{\nu}^{T_{i+1}, L}$ contain random terms. It simplifies the numerical implementation if all driving processes are non-homogeneous Lévy processes. This can be achieved by introducing an approximation, where we replace the random terms by their deterministic initial values at time $t=0$.

## 4 Pricing of swaptions

In this section we consider the pricing of swap rate based options, or swaptions, more precisely we shall price a swaption where the underlying swap starts at tenor time point $T_{i}$ and ends at time point $T_{N}$. By using general valuation results (see e.g. Musiela and Rutkowski (1997b)), the time $t=0$ price of the forward payer swaption is given by

$$
\begin{align*}
\mathrm{PS}_{0} & =\sum_{k=i+1}^{N} \delta_{k} B\left(0, T_{k}\right) \mathrm{E}^{\widetilde{\mathbb{P}}_{T_{i+1}}}\left[\left(S\left(T_{i}, T_{i}, T_{N}\right)-K\right)_{+}\right] \\
& =C_{i, N}(0) \mathrm{E}^{\widetilde{\mathbb{P}}_{T_{i+1}}}\left[\left(S\left(T_{i}, T_{i}, T_{N}\right)-K\right)_{+}\right] . \tag{51}
\end{align*}
$$

Raible (2000) proposed a method for the evaluation of European stock options in a Lévy setting by using bilateral (or, two-sided) Laplace transforms. The approach is based on the observation that the pricing formula for European options can be represented as a convolution. Whence one can use the fact that the bilateral Laplace transform of a convolution is the product of the bilateral Laplace transforms of the factors (the latter transforms are usually known explicitly). Inversion of the bilateral Laplace transform then yields the option prices as a function of the current price of the
underlying asset, and can be accomplished through the Fast Fourier Transform algorithm. See also the related work of Carr and Madan (1999), and Lee (2004).

Eberlein and Özkan (2005) showed how caplets can be priced using bilateral Laplace transforms. This method could also be employed for pricing the forward payer swaptions (51) as we shall shortly explain in the following. We do not give all technical assumptions or details accompanying that approach, instead we refer the reader to Raible (2000).

We concentrate on the purely discontinuous case ( $c=0$ in (32)) since we will use generalized hyperbolic Lévy processes in the implementations. We consider the forward swap rate in the form as given in (31):

$$
S\left(t, T_{i}, T_{N}\right)=S\left(0, T_{i}, T_{N}\right) \exp \left(\int_{0}^{t} \gamma\left(s, T_{i}\right) d \widetilde{L}_{s}^{T_{i+1}}\right)
$$

where the condition (33) now simplifies to

$$
\int_{0}^{t} \gamma\left(s, T_{i}\right) b_{s}^{T_{i+1}} d s=-\int_{0}^{t} \int_{\mathbb{R}}\left(\mathrm{e}^{\gamma\left(s, T_{i}\right) x}-1-\gamma\left(s, T_{i}\right) x\right) \widetilde{\nu}^{T_{i+1}, L}(d s, d x)
$$

We define

$$
X_{t}:=\int_{0}^{t} \gamma\left(s, T_{i}\right) d \widetilde{L}_{s}^{T_{i+1}}
$$

so that

$$
\begin{aligned}
X_{T_{i}} & =\ln \left(\frac{S\left(T_{i}, T_{i}, T_{N}\right)}{S\left(0, T_{i}, T_{N}\right)}\right) \\
& =\int_{0}^{T_{i}} \gamma\left(s, T_{i}\right) b_{s}^{T_{i+1}} d s+\int_{0}^{T_{i}} \int_{\mathbb{R}} x \gamma\left(s, T_{i}\right)\left(\mu^{L}-\widetilde{\nu}^{T_{i+1}, L}\right)(d s, d x)
\end{aligned}
$$

By defining $w(x, K):=(x-K)_{+}$, the payoff of the swaption is given by $w\left(S\left(T_{i}, T_{i}, T_{N}\right), K\right)$ and its price at time $t=0$ by $C_{i, N}(0) \mathrm{E}^{\widetilde{\mathbb{P}}_{T_{i+1}}}\left[w\left(S\left(T_{i}, T_{i}, T_{N}\right), K\right)\right]$. We consider the modified payoff $\widetilde{w}(x, K):=w\left(\mathrm{e}^{-x}, K\right)$. Let $\zeta_{i}:=-\ln S\left(0, T_{i}, T_{N}\right)$, then $S\left(T_{i}, T_{i}, T_{N}\right)=\mathrm{e}^{-\zeta_{i}+X_{T_{i}}}$. Furthermore, denote by $V\left(\zeta_{i}, K\right)$ the time zero price of the swaption, and let $L[\widetilde{w}]$ be the bilateral Laplace transform of $\widetilde{w}$ :

$$
L[\widetilde{w}](z)=\int_{-\infty}^{+\infty} \mathrm{e}^{-z x} \widetilde{w}(x) d x, \quad z=R+\mathrm{i} u \in \mathbb{C}, \quad R, u \in \mathbb{R}
$$

The price of the swaption at time zero can be written (apart from the discount factor) as a convolution of functions $\widetilde{w}(x)$ and $\rho(x)$, taken at the point $\zeta_{i}$ :

$$
\begin{aligned}
V\left(\zeta_{i}, K\right) & =C_{i, N}(0) \mathrm{E}^{\widetilde{\mathbb{P}}_{T_{i+1}}}\left[w\left(\mathrm{e}^{-\zeta_{i}+X_{T_{i}}}, K\right)\right] \\
& =C_{i, N}(0) \mathrm{E}^{\widetilde{\mathbb{P}}_{T_{i+1}}}\left[\widetilde{w}\left(\zeta_{i}-X_{T_{i}}, K\right)\right] \\
& =C_{i, N}(0) \int_{\mathbb{R}} \widetilde{w}\left(\zeta_{i}-x, K\right) \rho(x)(d x),
\end{aligned}
$$

where $\rho$ is the density of $X_{T_{i}}$. As mentioned above, the bilateral Laplace transform of a convolution equals the product of the bilateral Laplace transforms of the factors. Thus, we have that

$$
\begin{equation*}
L[V](R+\mathrm{i} u)=C_{i, N}(0) L[\widetilde{w}](R+\mathrm{i} u) \cdot L[\rho](R+\mathrm{i} u), \quad(u \in \mathbb{R}) \tag{52}
\end{equation*}
$$

By Theorem B. 2 in Raible (2000) the bilateral Laplace integral defining $L[V](z)$ converges absolutely, and $\zeta \mapsto V(\zeta, K)$ is a continuous function. Hence, according to Theorem B. 3 in Raible (2000) we can invert the bilateral Laplace transform to obtain the swaption price $V$ :

$$
\begin{align*}
V\left(\zeta_{i}, K\right) & =\frac{1}{2 \pi \mathrm{i}} \int_{R-\mathrm{i} \infty}^{R+\mathrm{i} \infty} \mathrm{e}^{\zeta_{i} z} L[V](z) d z \\
& =\frac{1}{2 \pi} \int_{-\infty}^{+\infty} \mathrm{e}^{\zeta_{i}(R+\mathrm{i} u)} L[V](R+\mathrm{i} u) d u \\
& =\frac{\mathrm{e}^{\zeta_{i} R}}{2 \pi} \lim _{M \rightarrow \infty}^{M \rightarrow \infty} \int_{-M}^{N} \mathrm{e}^{\mathrm{i} u \zeta_{i}} L[V](R+\mathrm{i} u) d u . \tag{53}
\end{align*}
$$

Note the identity $L[\rho](R+\mathrm{i} u)=\chi(\mathrm{i} R-u)$, where $\chi(\mathrm{i} R-u):=\mathrm{E}^{\tilde{\mathbb{P}}_{T_{i+1}}}\left[\mathrm{e}^{\mathrm{i}(\mathrm{i} R-u) X_{T_{i}}}\right]$ is the extended characteristic function of $X_{T_{i}}$. By substituting (52) into (53) we finally obtain the swaption pricing formula

$$
\begin{equation*}
V\left(\zeta_{i}, K\right)=C_{i, N}(0) \frac{\mathrm{e}^{\zeta_{i} R}}{2 \pi} \lim _{\substack{M \rightarrow \infty \\ N \rightarrow \infty}} \int_{-M}^{N} \mathrm{e}^{\mathrm{i} u \zeta_{i}} L[\widetilde{w}](R+\mathrm{i} u) \chi(\mathrm{i} R-u) d u . \tag{54}
\end{equation*}
$$

According to Raible (2000) it is sufficient to consider the case where the strike price equals one, since

$$
V\left(\zeta_{i}, K\right)=K V\left(\zeta_{i}+\ln K, 1\right) .
$$

The bilateral Laplace transform $L[\widetilde{w}]$ for $K=1$ is given by $L[\widetilde{w}](z)=(z(z+1))^{-1}$, if $\operatorname{Re} z<-1$. We remark that this approach can be used also for more complicated payoff functions as long as the payoff depends only on $X_{T_{i}}$ (for examples, we refer to Raible (2000)).

The characteristic function of $X_{T_{i}}$ has the following form and simplifies in the purely discontinuous case (see Eberlein and Raible (1999)):

$$
\begin{align*}
\chi(u):= & \mathrm{E}^{\widetilde{\mathbb{P}}_{T_{i+1}}}\left[\mathrm{e}^{\mathrm{i} u X_{T_{i}}}\right] \\
= & \exp \left(\mathrm{i} u \int_{0}^{T_{i}} \gamma\left(s, T_{i}\right) b_{s}^{T_{i+1}} d s\right. \\
& \left.\quad+\int_{0}^{T_{i}} \int_{\mathbb{R}}\left(\mathrm{e}^{\mathrm{i} u \gamma\left(s, T_{i}\right) x}-1-\mathrm{i} u \gamma\left(s, T_{i}\right) x\right) \widetilde{\nu}^{T_{i+1}, L}(d s, d x)\right) \\
= & \exp \left(-\mathrm{i} u \int_{0}^{T_{i}} \int_{\mathbb{R}}\left(\mathrm{e}^{\gamma\left(s, T_{i}\right) x}-1-\gamma\left(s, T_{i}\right) x\right) \widetilde{\nu}^{T_{i+1}, L}(d s, d x)\right. \\
& \left.\quad+\int_{0}^{T_{i}} \int_{\mathbb{R}}\left(\mathrm{e}^{\mathrm{i} u \gamma\left(s, T_{i}\right) x}-1-\mathrm{i} u \gamma\left(s, T_{i}\right) x\right) \widetilde{\nu}^{T_{i+1}, L}(d s, d x)\right) \\
= & \exp \left(\int_{0}^{T_{i}} \int_{\mathbb{R}}\left(\mathrm{e}^{\mathrm{i} u \gamma\left(s, T_{i}\right) x}-\mathrm{i} u \mathrm{e}^{\gamma\left(s, T_{i}\right) x}-(1-\mathrm{i} u)\right) \widetilde{\nu}^{T_{i+1}, L}(d s, d x)\right) . \tag{55}
\end{align*}
$$

The characteristic function (55) can be determined more precisely once the distribution of $L_{1}$ is specified (generalized hyperbolic or normal inverse Gaussian for instance). See e.g. Schoutens (2003) for a number of Lévy processes and their characteristic functions. Hence, we can calculate equation (54) numerically in an efficient way.

## 5 Concluding remarks

In this paper we derived the dynamics of the forward swap rate process in a semimartingale setting in explicit form and introduced a Lévy swap market model where swap rates are modeled as ordinary exponentials. Starting from the most distant swap rate (which is driven by a non-homogeneous Lévy process), we constructed the other swap rates such that they become martingales under the corresponding forward swap measures. We also showed how swaptions can be priced in this framework by using Laplace transform methods. A topic for future research will be to investigate further swap derivatives in the Lévy framework and to test the performance of the model.

## A Appendix

In this Appendix we will prove a general result which is applied in Section 3. The lemma corresponds to Lemma 2.3 in Rutkowski (2001), which we generalized to a semimartingale setting. It is also modified slightly to include the measure transform.

Lemma A. 1 Let $G$, H be real-valued adapted processes under some probability measure $\mathbb{P}$, satisfying the following SDEs

$$
\begin{align*}
& d G_{t}=g_{1}(t) d W_{t}+\int_{\mathbb{R}} g_{2}(t, x)(\mu-\nu)(d t, d x)  \tag{56}\\
& d H_{t}=h_{1}(t) d W_{t}+\int_{\mathbb{R}} h_{2}(t, x)(\mu-\nu)(d t, d x) \tag{57}
\end{align*}
$$

where $W_{t}$ is $\mathbb{P}$-Brownian motion and $\nu(d t, d x)$ is the $\mathbb{P}$-compensator of the random measure of jumps $\mu$. Let $g_{1}$, $h_{1}$ be square-integrable $\mathbb{P}$-a.s. and $g_{2}, h_{2} \in G_{\mathrm{loc}}(\mu)$. Suppose $H_{t}>-1$. Define $Y_{t}:=\left(1+H_{t}\right)^{-1}$. Then the process $Y G$ has the local martingale dynamics

$$
\begin{aligned}
d(Y G)_{t} & =Y_{t-}\left(g_{1}(t)-Y_{t-} G_{t-} h_{1}(t)\right) d \widetilde{W}_{t} \\
& +\int_{\mathbb{R}}\left(\frac{G_{t-}+g_{2}(t, x)}{1+H_{t-}+h_{2}(t, x)}-\frac{G_{t-}}{1+H_{t-}}\right)(\mu-\widetilde{\nu})(d t, d x)
\end{aligned}
$$

under a new measure $\widetilde{\mathbb{P}}, \widetilde{\mathbb{P}} \stackrel{\text { loc }}{<} \mathbb{P}$, where $\widetilde{W}_{t}$ is a $\widetilde{\mathbb{P}}$-Brownian motion, $d \widetilde{W}_{t}=d W_{t}-Y_{t-} h_{1}(t) d t$, and $\widetilde{\nu}(d t, d x)$ is the $\widetilde{\mathbb{P}}$-compensator of $\mu$ given by

$$
\widetilde{\nu}(d t, d x)=\left(1+Y_{t-} h_{2}(t, x)\right) \nu(d t, d x) .
$$

The density process is given by $\frac{d \widetilde{\mathbb{P}}}{d \mathbb{P}}=\mathcal{E}(M)$, where

$$
M_{t}=\int_{0}^{t} Y_{s-} h_{1}(s) d W_{s}+\int_{0}^{t} \int_{\mathbb{R}} Y_{s-} h_{2}(s, x)(\mu-\nu)(d s, d x) .
$$

Proof. We use the following short-hand notation in order to improve the readability: $g_{1}:=g_{1}(t)$, $h_{1}:=h_{1}(t), g_{2}:=g_{2}(t, x)$ and $h_{2}:=h_{2}(t, x)$. By applying Itô's formula for semimartingales (Jacod and Shiryaev (1987), I.4.57), we obtain

$$
d\left(\frac{G_{t}}{1+H_{t}}\right)=\frac{1}{1+H_{t-}} d G_{t}-\frac{G_{t-}}{\left(1+H_{t-}\right)^{2}} d H_{t}
$$

$$
\begin{align*}
& +\frac{1}{2}\left(-\frac{g_{1} h_{1}}{\left(1+H_{t-}\right)^{2}} d t+\frac{2 G_{t-} h_{1}^{2}}{\left(1+H_{t-}\right)^{3}} d t-\frac{g_{1} h_{1}}{\left(1+H_{t-}\right)^{2}} d t\right) \\
& +\int_{\mathbb{R}}\left(\frac{G_{t-}+g_{2}}{1+H_{t-}+h_{2}}-\frac{G_{t-}}{1+H_{t-}}-\left[\frac{g_{2}}{1+H_{t-}}-\frac{G_{t-} h_{2}}{\left(1+H_{t-}\right)^{2}}\right]\right) \mu(d t, d x) \\
= & \left(-\frac{g_{1} h_{1}}{\left(1+H_{t-}\right)^{2}}+\frac{G_{t-} h_{1}^{2}}{\left(1+H_{t-}\right)^{3}}\right) d t \\
& +\left(\frac{g_{1}}{1+H_{t-}}-\frac{G_{t-} h_{1}}{\left(1+H_{t-}\right)^{2}}\right) d W_{t} \\
& +\int_{\mathbb{R}} \frac{g_{2}}{1+H_{t-}}(\mu-\nu)(d t, d x)-\int_{\mathbb{R}} \frac{G_{t-} h_{2}}{\left(1+H_{t-}\right)^{2}}(\mu-\nu)(d t, d x) \\
& +\int_{\mathbb{R}}\left(\frac{G_{t-}+g_{2}}{1+H_{t-}+h_{2}}-\frac{G_{t-}}{1+H_{t-}}-\left[\frac{g_{2}}{1+H_{t-}}-\frac{G_{t-} h_{2}}{\left(1+H_{t-}^{2}\right)^{2}}\right]\right) \mu(d t, d x) . \tag{58}
\end{align*}
$$

The continuous part of (58) can be dealt with exactly the same way as in Lemma 2.3 in Rutkowski (2001), yielding

$$
\begin{align*}
& \left(-\frac{g_{1} h_{1}}{\left(1+H_{t-}\right)^{2}}+\frac{G_{t-} h_{1}^{2}}{\left(1+H_{t-}\right)^{3}}\right) d t+\left(\frac{g_{1}}{1+H_{t-}}-\frac{G_{t-} h_{1}}{\left(1+H_{t-}\right)^{2}}\right) d W_{t} \\
& =\left(Y_{t-}^{3} G_{t-} h_{1}^{2}-Y_{t-}^{2} g_{1} h_{1}\right) d t+\left(Y_{t-} g_{1}-Y_{t-}^{2} G_{t-} h_{1}\right) d W_{t} \\
& =Y_{t-}\left(g_{1}(t)-Y_{t-} G_{t-} h_{1}(t)\right)\left(d W_{t}-Y_{t-} h_{1}(t) d t\right) . \tag{59}
\end{align*}
$$

For the terms related to the jumps of $d\left(\frac{G_{t}}{1+H_{t}}\right)$ in (58), we obtain

$$
\begin{aligned}
& \int_{\mathbb{R}}\left(\frac{g_{2}}{1+H_{t-}}-\frac{G_{t-} h_{2}}{\left(1+H_{t-}\right)^{2}}\right)(\mu-\nu)(d t, d x) \\
& +\int_{\mathbb{R}}\left(\frac{G_{t-}+g_{2}}{1+H_{t-}+h_{2}}-\frac{G_{t-}}{1+H_{t-}}-\left[\frac{g_{2}}{1+H_{t-}}-\frac{G_{t-} h_{2}}{\left(1+H_{t-}\right)^{2}}\right]\right)(\mu-\nu)(d t, d x) \\
& +\int_{\mathbb{R}}\left(\frac{G_{t-}+g_{2}}{1+H_{t-}+h_{2}}-\frac{G_{t-}}{1+H_{t-}}-\left[\frac{g_{2}}{1+H_{t-}}-\frac{G_{t-} h_{2}}{\left(1+H_{t-}\right)^{2}}\right]\right) \nu(d t, d x) \\
& =\int_{\mathbb{R}}\left(\frac{G_{t-}+g_{2}}{1+H_{t-}+h_{2}}-\frac{G_{t-}}{1+H_{t-}}\right)(\mu-\nu)(d t, d x) \\
& \quad+\int_{\mathbb{R}}\left(\frac{G_{t-}+g_{2}}{1+H_{t-}+h_{2}}-\frac{G_{t-}}{1+H_{t-}}-\left[\frac{g_{2}}{1+H_{t-}}-\frac{G_{t-} h_{2}}{\left(1+H_{t-}\right)^{2}}\right]\right) \nu(d t, d x) .
\end{aligned}
$$

In view of Girsanov's theorem we can write this as

$$
\begin{align*}
& \int_{\mathbb{R}}\left(\frac{G_{t-}+g_{2}}{1+H_{t-}+h_{2}}-\frac{G_{t-}}{1+H_{t-}}\right)(\mu-\widetilde{\nu})(d t, d x) \\
& +\int_{\mathbb{R}}\left(\frac{G_{t-}+g_{2}}{1+H_{t-}+h_{2}}-\frac{G_{t-}}{1+H_{t-}}\right) \widetilde{Y}(t, x) \nu(d t, d x) \\
& +\int_{\mathbb{R}}\left(\frac{G_{t-} h_{2}}{\left(1+H_{t-}\right)^{2}}-\frac{g_{2}}{1+H_{t-}}\right) \nu(d t, d x), \tag{60}
\end{align*}
$$

where $\widetilde{\nu}(d t, d x)$ is a compensator of $\mu$ under some probability measure $\widetilde{\mathbb{P}} \stackrel{\text { loc }}{<} \mathbb{P}$. From the last equation it is evident that for $\frac{G_{t}}{1+H_{t}}$ to be a local martingale we must have

$$
\begin{align*}
& \int_{\mathbb{R}}\left(\frac{G_{t-}+g_{2}}{1+H_{t-}+h_{2}}-\frac{G_{t-}}{1+H_{t-}}\right) \tilde{Y}(t, x) \nu(d t, d x) \\
& +\int_{\mathbb{R}}\left(\frac{G_{t-} h_{2}}{\left(1+H_{t-}\right)^{2}}-\frac{g_{2}}{1+H_{t-}}\right) \nu(d t, d x)=0 . \tag{61}
\end{align*}
$$

From this we can determine the function $\widetilde{Y}(t, x)$ needed for the Girsanov change of measure:

$$
\begin{equation*}
\tilde{Y}(t, x)=\frac{1+H_{t-}+h_{2}}{1+H_{t-}} . \tag{62}
\end{equation*}
$$

The density process follows from Jacod and Shiryaev (1987, Theorem III.5.19).

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[^1]:    ${ }^{2}$ This is the measure associated to the numéraire $B\left(t, T_{N}\right)$.

[^2]:    ${ }^{3}$ For definition of this set we refer to Jacod and Shiryaev (1987, II.1.27)

